ON LINEAR AND NONLINEAR FRACTIONAL PDEs

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On linear and nonlinear fractional PDEs

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Abstract

In this study, Variational Iteration Method (VIM) has been applied to obtain the analytical solutions of fractional order nonlinear partial differential equations. The iteration procedure is based on a relatively new approach which is called Jumarie's fractional derivative. Several examples have been solved to elucidate effectiveness of the proposed method and the results are compared with the exact solution, revealing high accuracy and efficiency of the method.

Keywords: Fractional Fokker-Plank Equation, Fractional Fornberg-Whitham equation, Variational Iteration Method, Modified Riemann– Liouville derivative.

1 Introduction

In last two decades, fractional differential equations have gained much interest due to exact description of nonlinear phenomena in fluid flow, seismology, biology, chemistry, economic, probability and statistics, acoustics, material science, engineering and other areas of science. However, fractional calculus is three centuries old as the conventional calculus [1]. Derivatives and integrals of fractional arbitrary orders have found many applications in recent

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study [2,3]. Several analytical methods have been proposed for solving fractional differential equation of complex physical nature, such as Adomian's Decomposition Method (ADM)[4], Differential Transform Method (DTM) [5], Homotopy Perturbation Method (HPM) [6].

Obidat and Momani applied the VIM to fractional differential equations in fluid mechanics [7]. Inc [8] applied the VIM to solve the space and time fractional Burger's equations. Safari et al. [9] used the VIM to obtain an analytic approximate solution of space-fractional KdV-Burgers- Kuramoto equations. Abidi and Omrani [10] used the Homotopy Analysis Method (HAM) for the Fornberg-Whitham equation and Hassan et al. [19] applied (HAM) for solving Space and Time-fractional KdV Equations. Gupta and singh [11] applied the Homotopy Perturbation Method (HPM) for the fractional Fornberg-Whitham equation. More recently, Lu [12] has used the Variational Iteration Method (VIM) to obtain an approximate solution of Fornberg-Whitham equation. In this paper, we extend the application of the FVIM with modified Riemann-Liouville derivative in order derive the analytical approximate solutions to nonlinear time-fractional Fokker-Plank and Fornberg-Whitham equations.

2 Basic definitions

We give some basic definitions, notations and properties of the fractional calculus theory which are used further in this paper:

Definition 2.1 Assume $f : R \to R, x \to f(x)$, denote a continuous (but not necessarily differentiable) function and let the partition h > 0 in the interval [0, 1]. Jumarie's derivative is defined through the fractional difference:

$$\Delta^{\alpha} = (FW - 1)^{\alpha} f(x)$$

$$= \sum_{k=0}^{\infty} (-1)^{k} {\alpha \choose k} f(x + (\alpha - k)),$$
(1)

where, FWf(x) = f(x + h). Then the fractional derivative is defined as the following limit:

$$f^{(\alpha)} = \lim_{h \to 0} \frac{\Delta^{\alpha} \left[f(x) - f(0) \right]}{h^{\alpha}}.$$
 (2)

This definition is close to the standard definition of derivative, and as a direct result, the αth derivative of a constant $0 < \alpha < 1$; is zero.

Definition 2.2 The Riemann–Liouville fractional integral operator of order $\alpha \geq 0$ for a function $f \in C_{\mu} \mu \geq -1$ is defined as

$$I_x^{\alpha} f(x) = \frac{1}{\Gamma(\alpha)} \int_0^x (x - \xi)^{\alpha - 1} f(\xi) \, d\xi, \, \alpha > 0, \, t > 0.$$
(3)

Definition 2.3 The Jumarie's modified Riemann–Liouville derivative is defined as

$$I_x^{\alpha} f(x) = \frac{1}{\Gamma(m-\alpha)} \frac{d^m}{d x^m} \int_0^x (x-\xi)^{m-\alpha} (f(\xi) - f(0)) d\xi, \qquad (4)$$

where, $x \in [0, 1], m - 1 < \alpha \le m, m \ge 1$.

The proposed modified Riemann–Liouville derivative as shown in Eq. (4) is strictly equivalent to Eq. (2). Meanwhile, we would introduce some properties of the fractional modified Riemann–Liouville derivative

(a) Fractional Leibnitz product law

$$D_x^{\alpha}(u\,v) \,=\, u^{(\alpha)}\,v \,+\, u\,v^{(\alpha)}. \tag{5}$$

(b) Fractional Leibnitz formulation

$$I_x^{\alpha} D_x^{\alpha} f(x) = f(x) - f(0), \ 0 < \alpha \le 1.$$
(6)

Therefore, the integration by part can be used during the fractional calculus

$$I_x^{\alpha} u^{(\alpha)} v = (u v)|_a^b - I_b^{\alpha} u v^{(\alpha)}.$$
(7)

Definition 2.4 Fractional derivative of compounded functions is defined as

$$d^{\alpha}f(x) \cong \Gamma(1+\alpha) df, \ 0 < \alpha < 1.$$
(8)

Definition 2.5 The integral with respect to $(d\xi)^{\alpha}$ is defined as the solution of fractional differential equation given by equation

$$dy \cong f(x) (dx)^{\alpha}, y(0) = 0, x \ge 0, 0 < \alpha \le 1,$$
 (9)

$$y \cong \int_{0}^{x} f(\xi) (d\xi)^{\alpha} = \alpha \int_{0}^{x} (x - \xi)^{\alpha - 1} f(\xi) d\xi, \ 0 < \alpha \le 1.$$
(10)

For example $f(x) = x^{\beta}$ in Eq. (10), one obtains

$$\int_{0}^{x} \xi^{\beta} (d\xi)^{\alpha} = \frac{\Gamma(1+\alpha) \Gamma(1+\beta)}{\Gamma(1+\alpha+\beta)} x^{\alpha+\beta}, \ 0 < \alpha \le 1$$
(11)

and for a continuous function f, we have

$$\int_{0}^{x} \frac{\partial^{\alpha} f(t)}{\partial t^{\alpha}} \left(d\xi \right)^{\alpha} = \Gamma \left(1 + \alpha \right) \left(f(x) - f(0) \right).$$
(12)

Definition 2.6 Assume that the continuous function $f : R \to R, x \to f(x)$ has a fractional derivative of order $k\alpha$, for any positive integer k and any $\alpha 0 < \alpha \leq 1$, then the following equality holds, which is

$$f(x+h) = \sum_{k=0}^{\infty} \frac{h^{\alpha k}}{\alpha k!} f^{(\alpha k)}(x), \ 0 < \alpha \le 1,$$
(13)

On making the substitution $h \to x$ and $x \to 0$, we obtain the fractional Mc-Laurin series

$$f(x) = \sum_{k=0}^{\infty} \frac{x^{\alpha k}}{\alpha k!} f^{(\alpha k)}(0), \ 0 < \alpha \le 1.$$
 (13')

3 Analysis of Fractional Variational Iteration Method (FVIM)

In order to elucidate the solution procedure of the VIM, we consider the following fractional differential equation:

$$\frac{\partial^{\alpha}}{\partial t^{\alpha}} u\left(x, t\right) = K[x] u\left(x, t\right) + q\left(x, t\right), \ t > 0, \ x \in R,$$
(14)

subject to the initial condition

$$u(x,0) = f(x).$$

Here K[x] is the differential operator, f(x) and q(x, t) are continuous functions. According to VIM introduced by He [13], we can construct a correction functional for Eq. (14) as follows

$$u_{n+1}(x,t) = u_n(x,t) + I^{\alpha} \left[\lambda \left(\frac{\partial^{\alpha} u_n}{\partial t^{\alpha}} - K[x] u(x,t) - q(x,t) \right) \right],$$

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$$u_{n+1}(x,t) = u_n(x,t) + \frac{1}{\Gamma(\alpha)} \int_0^t (t-\xi)^{\alpha-1} \lambda(\xi) \Big(-q(x,\xi) + \frac{\partial^{\alpha} u_n}{\partial \xi^{\alpha}} (x,\xi) - K[x] u(x,\xi) \Big) d\xi.$$
(15)

Combining Eq.(10) and (15), we obtained a proposed correction functional

$$u_{n+1}(x,t) = u_n(x,t) + \frac{1}{\Gamma(\alpha+1)} \int_0^t \lambda\left(\xi\right) \left(\frac{\partial^{\alpha} u_n}{\partial \xi^{\alpha}}(x,\xi) - K\left[x\right] u\left(x,\xi\right) - q\left(x,\xi\right)\right) (d\xi)^{\alpha}.$$
(16)

It is obvious that the successive approximation u_j , $j \ge 0$ can be established by determining λ via variational theory. The function \tilde{u}_n is a restricted variation which means $\delta \tilde{u}_n = 0$. Therefore, we first determine Lagrange's multiplier that will be identified optimally via integration by parts. The successive approximation of the $u_{n+1}(x, t)$, $n \ge 0$ solution u(x, t) will be readily obtained upon using the Lagrange's multiplier and by using any selective function u_0 . The initial values are usually used for selecting the zeroth approximation u_0 . With λ determined, several approximations u_j , $j \ge 0$ follow immediately. Consequently, the exact solution may be obtained by using

$$u(x,t) = \lim_{n \to \infty} u_n(x,t) .$$
(17)

4 Numerical examples

In this section, we apply the FVIM to find the solutions of nonlinear timefractional Fokker-Plank and Fornberg-Whitham equations, and compared them with those obtained by other methods.

Example 1. Consider the following non-homogeneous space fractional equation

$$\frac{\partial^{\alpha} u}{\partial x^{\alpha}} - \frac{\partial u}{\partial x} = \frac{\partial u}{\partial t} + (2 - 2t - 2x), \quad (t > 0, \ x \in R, \ 1 < \alpha \le 2), \ (18)$$

with the conditions

$$u(0, t) = t^2, \quad u(x, 0) = x^2, \quad u_t(0, t) = 0.$$

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The correction functional is read as

$$u_{n+1}(x,t) = u_n(x,t) + \frac{1}{\Gamma(1+\alpha)} \int_0^x \lambda(\xi) \left(\frac{\partial^\alpha u_n(\xi,t)}{\partial \xi^\alpha} - \frac{\partial u_n(\xi,t)}{\partial \xi} - \frac{\partial u_n(\xi,t)}{\partial t} - (2-2t-2\xi)\right) (d\xi)^\alpha,$$
(19)

$$u_{n+1}(x,t) = u_n(x,t) + \frac{1}{\Gamma(1+\alpha)} \int_0^x \lambda(\xi) \left(\frac{\partial^\alpha u_n(\xi,t)}{\partial \xi^\alpha} - \frac{\partial \tilde{u}_n(\xi,t)}{\partial \xi} - \frac{\partial \tilde{u}_n(\xi,t)}{\partial t} - (2-2t-2\xi)\right) (d\xi)^\alpha.$$
(20)

The function \tilde{u}_n is a restricted term which means that $\delta \tilde{u}_n = 0$, via the variational theory. Taking variation δ to both sides of Eq. (19) and applying optimal condition $\delta u_{n+1} = 0$, one obtains

$$0 = \delta u_n + \frac{1}{\Gamma(1+\alpha)} \delta \int_0^x \lambda(\xi) \frac{\partial^\alpha u_n}{\partial \xi^\alpha} (d\xi)^\alpha, \qquad (21)$$

Using Eq. (8), applying integration by parts on second term of right hand and after equating, we have

$$\frac{\partial^{\alpha} \lambda\left(\xi\right)}{\partial \xi^{\alpha}} = 0, \text{ and } 1 + \lambda\left(\xi\right) = 0.$$
(22)

The generalized Lagrange multiplier can be identified by the above equations $\lambda(\xi) = -1$. Substituting Eq. (21) into the functional Eq. (19) yields the iteration formulation as follows

$$u_{n+1} = u_n - \frac{1}{\Gamma(1+\alpha)} \int_0^x \left(\frac{\partial^\alpha u_n}{\partial \xi^\alpha} - \frac{\partial u_n}{\partial \xi} - \frac{\partial u_n}{\partial t} - (2 - 2t - 2\xi) \right) (d\xi)^\alpha$$
(23)

The initial approximation is read as

$$u_0(x, t) = t^2 + \frac{2 - 2t}{\Gamma(1 + \alpha)} x^{\alpha} - \frac{2}{\Gamma(2 + \alpha)} x^{1 + \alpha}.$$

The other approximations are

$$u_{1}(x,t) = u_{0} - \frac{1}{\Gamma(1+\alpha)} \int_{0}^{x} \left\{ \frac{\partial^{\alpha} u_{0}}{\partial \xi^{\alpha}} - \frac{\partial u_{0}}{\partial \xi} - \frac{\partial u_{0}}{\partial t} \right\} (d\xi)^{\alpha}$$
$$= \frac{2-2t}{\Gamma(2\alpha)} x^{2\alpha-1} - \frac{4}{\Gamma(1+2\alpha)} x^{2\alpha} + \frac{2t}{\Gamma(1+\alpha)} x^{\alpha},$$
$$u_{2}(x,t) = u_{10} - \frac{1}{\Gamma(1+\alpha)} \int_{0}^{x} \left(\frac{\partial^{\alpha} u_{0}}{\partial t^{\alpha}} - \frac{\partial u_{0}}{\partial t^{\alpha}} - \frac{\partial u_{0}}{\partial t^{\alpha}} \right) (d\xi)^{\alpha}$$

$$\Gamma(1+\alpha) \int_{0} \left(\frac{\partial \xi^{\alpha}}{\partial \xi} - \frac{\partial \xi}{\partial t} \right)^{(\alpha\xi)} = \frac{2-2t}{\Gamma(3\alpha-1)} x^{3\alpha-2} - \frac{6}{\Gamma(3\alpha)} x^{3\alpha} + \frac{2t}{\Gamma(2\alpha)} x^{2\alpha-1} + \frac{2}{\Gamma(2+\alpha)} x^{2\alpha},$$

$$\vdots$$

The solution is

$$u(x,t) = t^{2} + \frac{2-2t}{\Gamma(1+\alpha)}x^{\alpha} - \frac{2}{\Gamma(2+\alpha)}x^{1+\alpha} + \frac{2-2t}{\Gamma(2\alpha)}x^{2\alpha-1} - \frac{4}{\Gamma(1+2\alpha)}x^{2\alpha} + \frac{2t}{\Gamma(1+\alpha)}x^{\alpha} + \frac{2-2t}{\Gamma(3\alpha-1)}x^{3\alpha-2} - \frac{6}{\Gamma(3\alpha)}x^{3\alpha} + \frac{2t}{\Gamma(2\alpha)}x^{2\alpha-1} + \frac{2}{\Gamma(2+\alpha)}x^{2\alpha} + \dots$$
(24)

which is exactly the same as the one that was obtained in [14] using the Adomian's decomposition method. Due to self canceling "noise" terms appearance, by replacing $\alpha = 2$ in Eq. (24), yields the exact solution:

$$u(x,t) = x^2 + t^2$$

Example 2. Consider the following nonlinear time-fractional hyperbolic equation [15, 16]

$$\frac{\partial^{\alpha} u}{\partial t^{\alpha}} = \frac{\partial}{\partial x} \left(u \frac{\partial u}{\partial x} \right), \qquad (t > 0, \ x \in R, \ 1 < \alpha \le 2), \qquad (25)$$

with the initial conditions

$$u(x, 0) = x^2, u_t(x, 0) = -2x^2.$$

The exact solution is $u(x, t) = \left(\frac{x}{t+1}\right)^2$. The correction functional is read as

$$u_{n+1} = u_n + \frac{1}{\Gamma(1+\alpha)} \int_0^t \lambda\left(\xi\right) \left(\frac{\partial^\alpha u_n}{\partial\xi^\alpha} - u_n \frac{\partial^2 u_n}{\partial x^2} - \left(\frac{\partial u_n}{\partial x}\right)^2\right) (d\xi)^\alpha$$
(26)

Making the correction functional stationary, the general Lagrange multiplier can be identified as $\lambda(\xi) = -1$. Substituting value of the general Lagrange multiplier in the Eq. (26), we get the following iteration formula

$$u_{n+1} = u_n - \frac{1}{\Gamma(1+\alpha)} \int_0^t \left(\frac{\partial^\alpha u_n}{\partial \xi^\alpha} - u_n \frac{\partial^2 u_n}{\partial x^2} - \left(\frac{\partial u_n}{\partial x}\right)^2\right) (d\xi)^\alpha \,.$$
(27)

The initial approximation is given as

$$u_0(x, t) = x^2 (1 - 2t),$$

The other approximations are

$$u_{1}(x,t) = u_{0} - \frac{1}{\Gamma(1+\alpha)} \int_{0}^{t} \left(\frac{\partial^{\alpha} u_{0}}{\partial \xi^{\alpha}} - u_{0} \frac{\partial^{2} u_{0}}{\partial x^{2}} - \left(\frac{\partial u_{0}}{\partial x} \right)^{2} \right) (d\xi)^{\alpha}$$
$$= \frac{6x^{2}}{\Gamma(1+\alpha)} \left(t^{\alpha} - \frac{4\Gamma(1+\alpha)}{\Gamma(2+\alpha)} t^{1+\alpha} + \frac{8\Gamma(1+\alpha)}{\Gamma(3+\alpha)} t^{2+\alpha} \right),$$
$$\vdots$$

Consequently, we have the following solution of Eq. (25) in a series form

$$u(x,t) = x^{2} (1-2t) + \frac{6x^{2}}{\Gamma(1+\alpha)} \left(t^{\alpha} - \frac{4\Gamma(1+\alpha)}{\Gamma(2+\alpha)} t^{1+\alpha} + \frac{8\Gamma(1+\alpha)}{\Gamma(3+\alpha)} t^{2+\alpha} \right) + \dots$$
(28)

Example 3. Consider the nonlinear time fractional Fokker - Plank equation

$$\frac{\partial^{\alpha} u}{\partial t^{\alpha}} = \left(-\frac{\partial}{\partial x}\left(\frac{4u}{x} - \frac{x}{3}\right) + \frac{\partial^2 u}{\partial x^2}\right) u\left(x, t\right), \quad (t > 0, x \in \mathbb{R}, 0 < \alpha \le 1),$$
(29)

subject to the initial condition

$$u\left(x,\,0\right)\,=\,x^2.$$

The correction functional is read as

$$u_{n+1} = u_n + \frac{1}{\Gamma(1+\alpha)} \int_0^t \lambda(\xi) \left(\frac{\partial^{\alpha} u_n}{\partial \xi^{\alpha}} + \frac{\partial}{\partial x} \left(\left(\frac{4u_n}{x} - \frac{x}{3}\right) u_n \right) - u_n \frac{\partial^2 u_n}{\partial x^2} \right) (d\xi)^{\alpha},$$
(30)

where, $\lambda\left(\xi\right)=-1$ can be determined optimally via variational theory. The iteration formula is

$$u_{n+1} = u_n - \frac{1}{\Gamma(1+\alpha)} \int_0^t \left(\frac{\partial^{\alpha} u_n}{\partial \xi^{\alpha}} + \frac{\partial}{\partial x} \left(\left(\frac{4u_n}{x} - \frac{x}{3}\right) u_n \right) - u_n \frac{\partial^2 u_n}{\partial x^2} \right) (d\xi)^{\alpha},$$
(31)

The initial approximation is given as

$$u_0\left(x,\,t\right)\,=\,x^2$$

$$\begin{split} u_1 &= u_0 - \frac{1}{\Gamma(1+\alpha)} \int_0^t \left(\frac{\partial^{\alpha} u_0}{\partial \xi^{\alpha}} + \frac{\partial}{\partial x} \left(\left(\frac{4u_0}{x} - \frac{x}{3} \right) u_0 \right) - u_0 \frac{\partial^2 u_0}{\partial x^2} \right) (d\xi)^{\alpha} \\ &= \frac{x^2}{\Gamma(1+\alpha)} t^{\alpha}, \\ u_2 &= u_1 - \frac{1}{\Gamma(1+\alpha)} \int_0^t \left(\frac{\partial^{\alpha} u_1}{\partial \xi^{\alpha}} + \frac{\partial}{\partial x} \left(\left(\frac{4u_1}{x} - \frac{x}{3} \right) u_1 \right) - \frac{\partial^2}{\partial x^2} \left(u_1^2 \right) \right) (d\xi)^{\alpha} \\ &= \frac{x^2}{\Gamma(1+2\alpha)} t^{2\alpha}, \end{split}$$

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$$u_{3} = u_{2} - \frac{1}{\Gamma(1+\alpha)} \int_{0}^{t} \left(\frac{\partial^{\alpha} u_{2}}{\partial \xi^{\alpha}} + \frac{\partial}{\partial x} \left(\left(\frac{4u_{2}}{x} - \frac{x}{3} \right) u_{2} \right) - u_{2} \frac{\partial^{2} u_{2}}{\partial x^{2}} \right) (d\xi)^{\alpha}$$
$$= \frac{x^{2}}{\Gamma(1+3\alpha)} t^{3\alpha},$$
$$\vdots$$

The series solution is given by

$$u(x,t) = x^{2} \left(1 + \frac{1}{\Gamma(1+\alpha)} t^{\alpha} + \frac{1}{\Gamma(1+2\alpha)} t^{2\alpha} + \frac{1}{\Gamma(1+3\alpha)} t^{3\alpha} \dots \right).$$
$$= \sum_{k=0}^{\infty} \frac{x^{2}}{\Gamma(1+k\alpha)} t^{k\alpha}$$
$$= x^{2} E_{\alpha}(t^{\alpha}).$$
(32)

which is the exact solution of the nonlinear time fractional PDE obtained by Odibat and Momani [17].

Example 4. Consider the following nonlinear time-fractional Fornberg-Whitham equation

$$\frac{\partial^{\alpha} u}{\partial t^{\alpha}} - \frac{\partial^{3} u}{\partial x^{2} \partial t} + \frac{\partial u}{\partial x} = u \frac{\partial^{3} u}{\partial x^{3}} - u \frac{\partial u}{\partial x} + 3 \frac{\partial u}{\partial x} \frac{\partial^{3} u}{\partial x^{3}}, \quad (t > 0, x \in R, 0 < \alpha \le 1),$$
(33)

with initial condition as

$$u(x, 0) = \frac{4}{3}e^{\frac{1}{2}x}.$$

The exact travelling wave solution to the above initial value problem is given by [18]

$$u(x, t) = \frac{4}{3}e^{\frac{1}{2}x - \frac{2}{3}t}.$$

The correction functional is read as

$$u_{n+1} = u_n + \frac{1}{\Gamma(1+\alpha)} \int_0^t \lambda\left(\xi\right) \left(\frac{\partial^{\alpha} u_n}{\partial \xi^{\alpha}} - \frac{\partial^3 u_n}{\partial x^2 \partial t} + \frac{\partial u_n}{\partial x} - u_n \frac{\partial^3 u_n}{\partial x^3} + u_n \frac{\partial u_n}{\partial x} - 3 \frac{\partial u_n}{\partial x} \frac{\partial^3 u_n}{\partial x^3}\right) (d\xi)^{\alpha}, \qquad (34)$$

Making the correction functional stationary, the general Lagrange multiplier can be identified as $\lambda(\xi) = -1$. Substituting value of the general Lagrange multiplier in the Eq. (34), we get the following iteration formula

$$u_{n+1} = u_n - \frac{1}{\Gamma(1+\alpha)} \int_0^t \left(\frac{\partial^{\alpha} u_n}{\partial \xi^{\alpha}} - \frac{\partial^3 u_n}{\partial x^2 \partial t} + \frac{\partial u_n}{\partial x} - u_n \frac{\partial^3 u_n}{\partial x^3} + u_n \frac{\partial u_n}{\partial x} - 3 \frac{\partial u_n}{\partial x} \frac{\partial^3 u_n}{\partial x^3} \right) (d\xi)^{\alpha}.$$
(35)

We start with an initial approximation

$$u_0(x, t) = \frac{4}{3}e^{\frac{1}{2}x}.$$

Then, we can obtain the other components by means of

$$\begin{split} u_1 &= u_0 - \frac{1}{\Gamma\left(1+\alpha\right)} \int\limits_0^t \left(\frac{\partial^{\alpha} u_0}{\partial \xi^{\alpha}} - \frac{\partial^3 u_0}{\partial x^2 \partial t} \right. \\ &+ \frac{\partial u_0}{\partial x} - u_0 \frac{\partial^3 u_0}{\partial x^3} + u_0 \frac{\partial u_0}{\partial x} - 3 \frac{\partial u_0}{\partial x} \frac{\partial^3 u_0}{\partial x^3} \right) (d\xi)^{\alpha} \\ &= -\frac{2}{3\Gamma\left(1+\alpha\right)} e^{\frac{1}{2}x} t^{\alpha}, \end{split}$$

$$\begin{split} u_2 &= u_1 - \frac{1}{\Gamma(1+\alpha)} \int_0^t \left(\frac{\partial^\alpha u_1}{\partial \xi^\alpha} - \frac{\partial^3 u_1}{\partial x^2 \partial t} \right. \\ &+ \frac{\partial u_1}{\partial x} - u_1 \frac{\partial^3 u_1}{\partial x^3} + u_1 \frac{\partial u_1}{\partial x} - 3 \frac{\partial u_1}{\partial x} \frac{\partial^3 u_1}{\partial x^3} \right) (d\xi)^\alpha \\ &= \frac{1}{6} e^{\frac{1}{2}x} \left(\frac{2}{\Gamma(1+\alpha)} t^{2\alpha} - \frac{1}{\Gamma(2\alpha)} t^{2\alpha-1} \right), \\ u_3 &= \frac{1}{72} e^{\frac{1}{2}x} \left(\frac{12}{\Gamma(3\alpha)} t^{3\alpha-1} - \frac{3}{\Gamma(3\alpha-1)} t^{3\alpha-2} \right. \\ &- \frac{12}{\Gamma(3\alpha+1)} t^{3\alpha} \right), \end{split}$$

and so on, in the same manner the remaining components can be obtained.

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The series solution is

$$u(x,t) = \frac{1}{72} e^{\frac{1}{2}x} \left(96 - \frac{48}{\Gamma(1+\alpha)} t^{\alpha} + \frac{24}{\Gamma(1+2\alpha)} t^{2\alpha} - \frac{12}{\Gamma(2\alpha)} t^{2\alpha-1} - \frac{12}{\Gamma(3\alpha)} t^{3\alpha-1} - \frac{3}{\Gamma(3\alpha-1)} t^{3\alpha-2} - \frac{12}{\Gamma(3\alpha+1)} t^{3\alpha} \dots \right).$$
(36)

5 Conclusion

In this study, Variational iteration method with new approach has been successfully employed to obtain exact and approximate analytical solutions of nonlinear fractional Fokker-Plank and Fornberg-Whitham equations. The method has been used in a direct way without linearization, perturbation or any restrictive assumption.

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O linearnim i nelinearnim frakcionim parcijalnim diferencijalnim jednačinama

Varijacioni iteracioni metod (VIM) je primenjen za dobijanje analitičkih rešenja nelinearnih frakcionih parcijalnih diferencijalnih jednačina. Iteracioni postupak se zasniva na relativno novom pristupu koji se zove Jumarie-ov frakcioni izvod. Nekoliko primera su rešeni u cilju razjašnjenja efikasnosti predloženog metoda i rezultati su uporedjeni sa tačnim rešenjima, otkrivajuci visoku tačnost i efikasnost metoda.

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